

Sterile Neutrinos in the MINSIS Experiment

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In collaboration with Jacobo López-Pavón

- Brief overview of MINSIS
- Sterile neutrinos: the $3 + 1$ scenario
- MINSIS sensitivity - dependence on:
 - Efficiency
 - Background
 - Beam energy
 - Flux
- Summary.

MINSIS = Main Injector Non-Standard Interactions Search

Experiment overview (see Adam Para's talk for details):

- Beam: NuMI beam \Rightarrow mainly ν_μ with some ν_e and ν_τ background.
- Baseline: $L = 1$ km.
- Detector: mass ~ 1 ton ($\sim 10^6$ CC ν_μ events), can detect τ 's
 \Rightarrow Look for $\nu_\mu \rightarrow \nu_\tau$ oscillations.

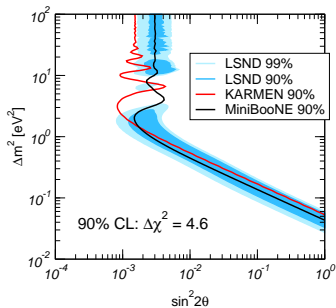
Sterile ν : the $3 + 1$ model

- The $3 + 1$ model is excluded at 4σ by a global fit to all available data.
M. Maltoni, T. Schwetz, [arXiv:0705.0107 [hep-ph]].
- But if we neglect only the LSND data, and keep data from all the remaining experiments, the model is no longer excluded.
- There are 3 additional mixing angles (θ_{41} , θ_{42} , θ_{43}) involved, and 2 additional CP violating phases.
- Different experiments put bounds on different combinations of these parameters.

Sterile ν : current bounds

- E.g. LSND, MiniBooNE, KARMEN, looked for $\nu_\mu \rightarrow \nu_e$ ($\bar{\nu}_\mu \rightarrow \bar{\nu}_e$) \Rightarrow bounds on $4 \sin^2 \theta_{14} \sin^2 \theta_{24} = \sin^2 2\theta$:

M. Maltoni, T. Schwetz, [arXiv:0705.0107 [hep-ph]].



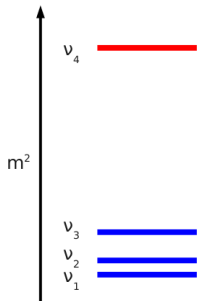
- MINSIS (like NOMAD and CHORUS) will be sensitive to $4 \cos^2 \theta_{14} \sin^2 \theta_{24} \cos^2 \theta_{24} \sin^2 \theta_{34} = \sin^2 2\theta_s^{\mu\tau}$ (no sensitivity to CP phases).

The 3 + 1 model: short baseline approximation

At $L = 1$ km, the **solar and atmospheric oscillations are irrelevant**:

$$\sin^2 \left(\frac{\Delta m_{21}^2 L}{4E} \right) \sim 10^{-10}$$

$$\sin^2 \left(\frac{\Delta m_{32}^2 L}{4E} \right) \sim 10^{-7}$$



So we can safely consider a **2 family approximation**, treating the 3 standard model ν as a single state.

In this approximation, we assume:

$$m_{\nu_4} \gg m_{\nu_3}, m_{\nu_2}, m_{\nu_1} \Rightarrow \Delta m_{43}^2 \sim \Delta m_{42}^2 \sim \Delta m_{41}^2 = \Delta m_s^2.$$

The 3 + 1 model: parameterization

The extended mixing matrix can be parameterized as:

J. N. Abdurashitov *et al.*, [arXiv:astro-ph/9907113].

$$U = R_{34}(\theta_{34})R_{24}(\theta_{24})R_{23}(\theta_{23}, \delta_3)R_{14}(\theta_{14})R_{13}(\theta_{13}, \delta_2)R_{12}(\theta_{12}, \delta_1)$$

where R_{ij} are 2×2 rotation matrices.

Using this parameterization and the SBL approximation, we obtain the following oscillation probabilities for a beam of ν_μ :

The 3 + 1 model: oscillation probabilities

$$P_{\nu_\mu \rightarrow \nu_e} = 4s_{14}^2 c_{14}^2 s_{24}^2 \sin^2 \Delta_s \quad (\text{e.g. MiniBooNE})$$

$$P_{\nu_\mu \rightarrow \nu_\mu} = 1 - 4c_{14}^2 s_{24}^2 (1 - c_{14}^2 s_{24}^2) \sin^2 \Delta_s \quad (\text{e.g. MINOS})$$

$$P_{\nu_\mu \rightarrow \nu_\tau} = 4c_{14}^4 s_{24}^2 c_{24}^2 s_{34}^2 \sin^2 \Delta_s \quad (\text{e.g. NOMAD})$$

$$P_{\nu_\mu \rightarrow \nu_s} = 4c_{14}^2 s_{24}^2 c_{24}^2 c_{34}^2 \sin^2 \Delta_s$$

where $s_{ij} = \sin \theta_{ij}$, $c_{ij} = \cos \theta_{ij}$ and $\Delta_s = \frac{\Delta m_s^2 L}{4E}$.

The experimental sensitivity will be affected by:

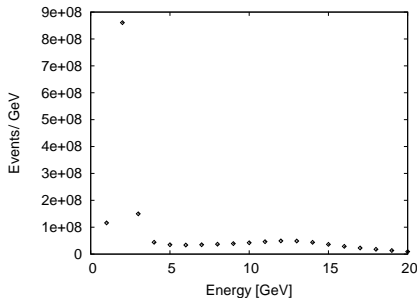
- Flux
- Energy
- Efficiency
- Backgrounds
- Systematics
- Energy resolution
- Anything else?

Beam: $\text{NO}\nu\text{A}$ flux

Initially, consider $\text{NO}\nu\text{A}$ flux, but only consider ν_μ for now:

D. Ayres *et al.*, [arXiv:hep-ex/0503053].

- 10^{21} PoT per year, 1.12 MW, running for 3 years.
- At $L = 1$ km, statistics are not a problem!
- Using a 4 kton detector, have $\sim 10^9$ CC ν_μ events for $1 \text{ GeV} < E < 20 \text{ GeV}$.
- Peak is at $\sim 2 \text{ GeV}$; have $\sim 10^8$ ν_μ above τ threshold.



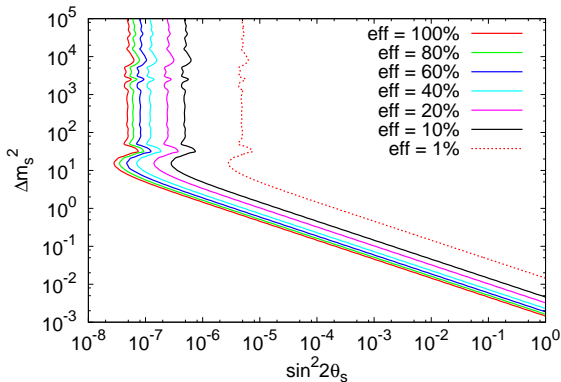
Experimental sensitivities

First simulate $\sin^2 2\theta_s^{\mu\tau} = \Delta m_s = 0$ (using GLoBES).

What bounds can we obtain on these parameters?

Will they be significantly better than current bounds?

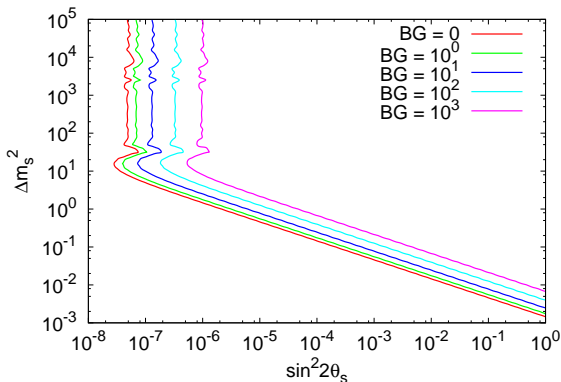
Consider a hypothetical **zero background**, and look at the effect of **altering the efficiency**:



Can place an upper bound on $\sin^2 2\theta_s^{\mu\tau}$ between $\sim 10^{-7}$ (100% efficiency) and $\sim 10^{-5}$ (1% efficiency) at 90% confidence.

Backgrounds

Consider a hypothetical **100% efficiency**, and look at the effect of **altering the number of background events**:



Can place an upper bound on $\sin^2 2\theta_s^{\mu\tau}$ between $\sim 10^{-7}$ (no background) and $\sim 10^{-6}$ (10^3 background events) at 90% confidence.

Is it worth detecting τ 's?

If we don't detect τ 's, then the other option is to detect μ 's
 \Rightarrow Look at μ disappearance channel, $\nu_\mu \rightarrow \nu_\mu$.

Can't compare $\nu_\mu \rightarrow \nu_\mu$ and $\nu_\mu \rightarrow \nu_\tau$ directly because these channels are sensitive to different combinations of angles:

$$P_{\nu_\mu \rightarrow \nu_\mu} = 1 - 4c_{14}^2 s_{24}^2 (1 - c_{14}^2 s_{24}^2) \sin^2 \Delta_s = 1 - \sin^2 2\theta_s^{\mu\mu} \sin^2 \Delta_s$$

$$P_{\nu_\mu \rightarrow \nu_\tau} = 4c_{14}^4 s_{24}^2 c_{24}^2 s_{34}^2 \sin^2 \Delta_s = \sin^2 2\theta_s^{\mu\tau} \sin^2 \Delta_s.$$

But can qualitatively compare sensitivities to $\sin^2 2\theta_s^{\mu\mu}$ and $\sin^2 2\theta_s^{\mu\tau}$.

τ appearance vs μ disappearance

Compare sensitivity to $\sin^2 2\theta_s^{\mu\mu}$ from 'perfect' μ detection (100% efficiency, zero background, 0.1% systematic errors)

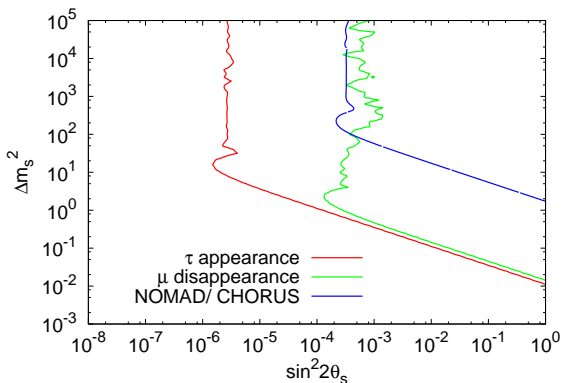
with sensitivity to $\sin^2 2\theta_s^{\mu\tau}$ from 'realistic' τ detection (10% efficiency, 2 background events, 10% systematic errors)

and sensitivity to $\sin^2 2\theta_s^{\mu\tau}$ from NOMAD and CHORUS.
E. Eskut *et al.*, [arXiv:0710.3361 [hep-ex]].

τ appearance vs μ disappearance

With μ disappearance, have similar sensitivity to $\sin^2 2\theta_5^{\mu\mu}$ as NOMAD and CHORUS have to $\sin^2 2\theta_5^{\mu\tau}$.

τ appearance is ~ 100 times more sensitive than μ disappearance.



Comparison with detector based on emulsion-silicon target

Yesterday we saw results from a 1 ton detector with 50% efficiency (2×10^6 CC ν_μ events detected):

J. J. Gomez-Cadenas, J. A. Hernando, Nuc. Inst. A, 281 (1996), 223-235.

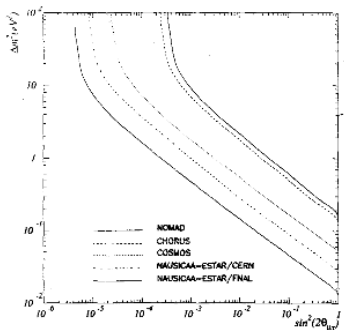


Fig. 11. Exclusion plot, showing the expected sensitivity of NAUSICAA-BSTAR at CERN and FNAL beams.

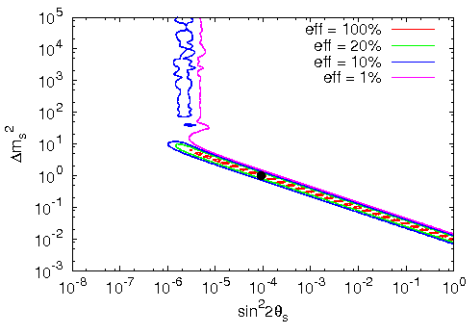
Detecting non-zero $\theta_s^{\mu\tau}$, Δm_s^2

Suppose now that $\sin^2 2\theta_s^{\mu\tau} = 10^{-4}$ and $\Delta m_s^2 = 1 \text{ eV}^2$.

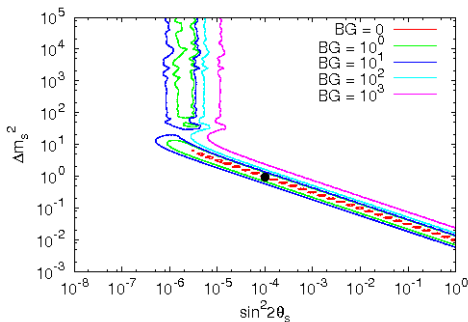
How precisely could we measure these parameters?

Obtain $\sim 100 \nu_\tau$ events in total (taking into account τ cross-section, and assuming 100% efficiency).

Zero background, vary efficiency

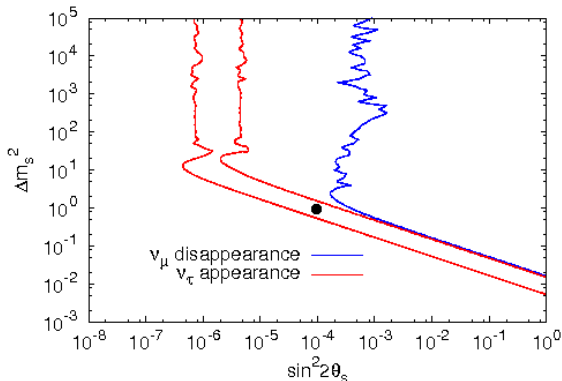


100% efficiency, vary background



Detecting non-zero $\theta_s^{\mu\tau}$, Δm_s^2

Again, compare 'perfect' μ detection with 'realistic' τ detection:

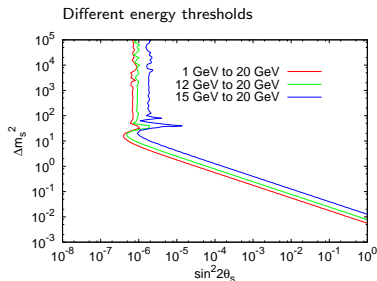
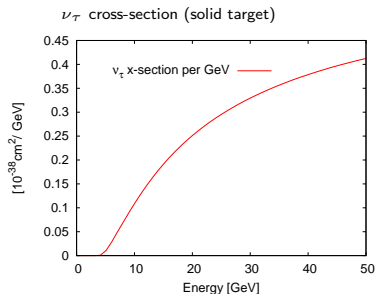


With 10% efficiency and 2 background events, MINSIS can significantly constrain the parameter space.

Energy dependence

The energy dependence of the sensitivities comes from the energy dependence of the ν_T cross-section.

M. Messier, UMI-99-23965; E. Paschos, J. Yu, [arXiv:hep-ph/0107261].

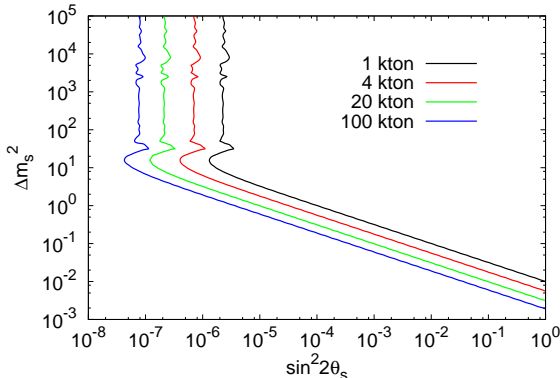


Events below ~ 12 GeV don't help much.

\Rightarrow **A higher energy beam, above ~ 12 GeV, would be good!**

Flux dependence

Can increase the number of higher energy events by increasing beam energy, running for longer, using a larger detector:



Roughly: **a 25 times increase in flux improves sensitivity by an order of magnitude.**

- Detecting $\nu_\mu \rightarrow \nu_\tau$ oscillations in the **MINSIS experiment could improve on current bounds on the 3+1 model of sterile neutrinos, given an adequately large detector/ enough events!**
- The critical factors are **efficiency, background rejection and energy.**
- Energy resolution and systematics have very little effect.